Test suite for NEIGHCRYS_2.3.0 and DMACRYS_2.3.0

This suite provides a set of files for testing new installations of **NEIGHCRYS_2.3.0** and **DMACRYS_2.3.0** and above. The comments, instructions and examples below are to supplement information given in the paper describing the DMACRYS programme¹ and the associated manual, to act as illustrations for users (but not necessarily of best practice, as designed for fast testing).

The directory test#_testing_REFCODE contains a README and directories including NEIGHCRYS_input, which contains the input files:

REFCODE.res (Crystal structure), bondlengths (define molecule bonds), dmacrys.dma (Distributed multipoles) dmacrys.mols (Molecular Axes) pote.dat/fit.pots/will01.pots (repulsion-dispersion model).

If you want to run NEIGHCRYS interactively, then the answers you need are in neighcrys answers.

Test#_molecule#	Example of & crystal structure	Molecular diagram
01_1	Lattice energy minimization with FIT exp-6 potential and distributed multipoles. The example files are for a computational model corresponding to AXOSOW.	\checkmark
02_1	Lattice energy minimization with WILLIAMS potential , where interaction sites on H are shifted, as contrast to 01. The example files are for a computational model corresponding to AXOSOW.	\checkmark
03_11	Lattice energy minimization with a custom potential . The example files are for a computational model corresponding to SOXLEX.	
04_11	Lattice energy minimization with a custom potential and without splines . The example files are for a computational model SOXLEX. The input files for NEIGHCRYS are identical to example 3, but the output file (SOXLEX_af395.dmain) has been edited before running DMACRYS .	
05_111	Lattice energy minimization of a large Z'>1 unit cell with 15 Å cutoff. The example files are from the CSD entry KAXXAI02.	J.J.

All the test examples were created on Xenon using the above versions of the codes.

06_111	Lattice energy minimization of a large Z'>1 unit cell with 15 Å and without splines. The example files are from the CSD entry KAXXAI02. The input files for NEIGHCRYS are identical to example 5, but the output file (KAXXAI02.dmain) has been edited before running DMACRYS.	J.t.
07_111	Lattice energy minimization of a large Z'>1 unit cell with 30 Å and without splines. The example files are from the CSD entry KAXXAI02. The input files for NEIGHCRYS are identical to example 5, but the output file (KAXXAI02.dmain) has been edited before running DMACRYS.	TTT T
08_IV	Lattice energy minimization of a small 2 component system. The example files are for a computational model corresponding to KONTIQ01.	
09_V	Lattice energy minimization of a crystal structure containing a spherical ion . The example files are for a computational model corresponding to FINVAZ.	,CI1
10_VI	Lattice energy minimization of a complex salt . The example files are for a computational model corresponding to WEMGEK.	Y H
11_VII	Properties calculation of a crystal structure. The example files are for a computational model corresponding to CBMZPN10. The NEIGHCRYS output file (CBMZPN_opt_III.dmain) has been edited before running DMACRYS .	

12_VIII	Properties calculation of a crystal structure using an anisotropic repulsion custom potential. The example files are for a computational model corresponding to DCLBEN06. The NEIGHCRYS output file (DCLBEN_beta_abi1.dmain) has been edited before running DMACRYS.	
13_V	Properties calculation of a crystal structure containing a spherical ion FINVAZ. The example files are for a computational model corresponding to FINVAZ. This is also a large cell (C2/c Z=8), so takes longer to run.	,CI1
14_IX	Lattice energy minimization with pressure . The example files are for the experimentally observed form at high pressure (0.5 GPa) of 4-fluorotoluene (a redetermination was published as YICDIZ01, although the pressure wasn't specified in the CSD entry). The NEIGHCRYS output file (YICDIZ.dmain) has been edited before running DMACRYS to specify the pressure.	
15_X	Lattice energy minimization with induction . The example files are for a computational model corresponding to OBEQUJ.	
16_XI	Symmetry reduction . The example files are for a computational model corresponding to PAPTUX, which corresponds to a transition state. The run which required this symmetry reduction is given for information in the directory reasons.	1 HA
17_XII	Pasting. The example files are for TEVSOD, with a molecule optimized using GAUSSIAN being substituted for the crystallographically determined molecular structure.	A A A A
18_1	Using default values. This is a run of NEIGHCRYS which uses the default values for filenames. The example files are for a computational model corresponding to AXOSOW.	\downarrow

19_XIII	A large cubic cell. This run of NEIGHCRYS has a Z'=1 two-component structure with 24 molecules in the whole cell. It failed with previous versions of the code.	A A A
20_XIII	Use of hessian. This run of DMACRYS uses the hessian output from a previous run of DMACRYS . This would normally be used to speed up a minimisation, either as a restart or where there is a small change in the starting structure or potential model, so that the Hessian from another run is a reasonable approximation.	A A A
21_XIV	Use of extended mode of NEIGHCRYS. This example has a long molecule, which the default settings of NEIGHCRYS cannot see all the atoms as connected. The extended mode of running NEIGHCRYS is required.	TH TH
22_XV	Use of damped dispersion. An example for use with non-empirical C_6 , C_8 and C_{10} dispersion models.	
23_111	Lattice energy minimization of a large Z'>1 unit cell with 30 Å and with splines. In the case of example 7, the NEIGHCRYS output file (KAXXAI02.dmain) had to be edited manually to double the cutoff distance; this example uses the extended mode of NEIGHCRYS to get a 30 Å cutoff.	
24_XVI	Use of non-empirical potentials. Illustrating the use of more complex anisotropic repulsion terms and damped dispersion.	
25_XVI	Energy refinement to include polarizability.	

I - AXOSOW 01.lem_fit_AXOSOW & 02.lem_will01_AXOSOW

This molecule was set as one of the challenges in the Fourth Blind Test of Crystal Structure Prediction.² The starting files are trial crystal structures generated by **MOLPAK**, with the distributed multipoles (dma) generated using **GAUSSIAN**³ and **GDMA**,⁴ as in all these examples.

The examples are given to show the differences in the FIT and Williams exp-6 potentials, whose coefficients are given in pote.dat (generic file type) file called fit.pots or will01.pots respectively. The two potentials do not only differ in the coefficients in pote.dat, but critically in the Williams potential has the hydrogen atom interaction sites moved in along the bonds by 0.1 Å to reflect the position of the centre of electron density. Further details of these two potentials, which are widely used for modelling organic crystals are given in the SI of the DMACRYS paper.¹ The wider range of atomic types and need for **NEIGHCRYS** to set up the shift the hydrogen interaction site positions for the Williams potential, leads to specific questions in the **NEIGHCRYS** input. The first difference in the output is the atomic types as seen in the fort.21 (**NEIGHCRYS** output), which are used to describe the atoms in the dmacrys.mols, dmacrys.dma and potential (fit.pots or will01.pots in these examples).

I_1 – FIT potential		I_2 – Williams potential	
Inequivalent basis atoms		Inequivalent basis atoms	
atom atomic name input index number name index 1 6 C F1 1 C1 2 6 C F1 2 2 3 6 C F1 3 3 C3 4 8 0 F1 1 4 01 5 1 H F1 1 5 H1 6 1 H F1 2 6 H2 7 1 H F1 3 7 H3 8 1 H F1 4 8 H4	molecule invert number flag 1 F 1 F 1 F 1 F 1 F 1 F 1 F 1 F 1 F	atom atomic name input molecule invert index number name name number flag 1 6 C_W3_1_1 Cl 1 F 2 6 C_W3_2 Cl 1 F 3 6 C_W3_3_3 C3 1 F 4 8 O_W1_14 O1 1 F 5 1 H_W1_16 H1 1 F 6 1 H_W1_26 H2 1 F 7 1 H_W1_37 H3 1 F 8 1 H_W1_48 H4 1 F	

The new hydrogen interactions sites required if Williams and foreshortening are requested in the **NEIGHCRYS** input are also given in fort.21

Hydrogen positions have been foreshortened. Positions of all atoms (with hydrogens not foreshortened) in the local axis system and centre of mass of the molecul- with foreshortened hydrogens for molecule 1	e
basis No. Species x y z (Angstroms) Mass	
1 C W3 1 -1.458147 -1.140153 0.000056 1 9 C W3 2 -0.777876 0.015814 0.000056 1 17 C W3 3 0.694209 0.015814 0.000056 1 25 O W1 1 1.379980 1.032476 -0.000193 1 33 H W1 1 -0.937054 -2.089416 0.001293 41 H W1 2 -2.537575 -1.172471 0.000375 49 H W1 3 -1.271827 0.978748 -0.001202 57 H W1 4 1.158373 -0.988909 0.000688	2.010700 2.010700 5.999400 1.007940 1.007940 1.007940 1.007940
Foreshortened hydrogen atom positions in the same local axis system	
basis No. Species x y z (Angstroms) Mass	
33 H_W1_1	1.007940 1.007940 1.007940 1.007940

This data is used to calculate the distributed multipoles with the H nuclei in the foreshortened positions, using options in GDMA,⁴ to analyse the ab initio wavefunction with the H nuclei in their true positions. Hence, the dma's in the two examples differ: not only in the H atom positions and multipoles, but also the non-hydrogen atoms will have the same positions (relative to the local axis system) but different multipoles because of the change in allocation of the electron density between the sites caused by the shift in H positions. However, as shown below, although the various contributions to the lattice energy change, the total electrostatic contribution to the initial lattice energy

does not change by much compared with the difference in the repulsion-dispersion potential. Note that the change in potential does make a considerable difference to the structure at the lattice energy minimum.

FIT First evaluation of lattice energy		Williams First evaluation of lattice energy		
Contributions to lattice energy (eV per unit	t cell [kJ/mol])	Contributions to lattice energy (eV per uni	t cell [kJ/mol])	
Z = 8	0 0024416051016	Z = 8	0 7600007001016	
-0.10883959E+03 1	-0.902441608+01[-0.92757803E+02 1	-0./69099/2E+01[
Intra-molecular charge-charge energy=	-0.78227800E+01[Intra-molecular charge-charge energy=	-0.67563847E+01[
-0.94347178E+02]		-0.81485844E+02]		
Inter-molecular charge-charge energy= -0.14492414E+02 1	-0.12016360E+01[Inter-molecular charge-charge energy= -0.11271959E+02 1	-0.93461253E+00[
Total charge-dipole energy=	0.20422449E+00[Total charge-dipole energy=	0.39674038E-02[
0.24630636E+01]		0.47849147E-01]		
Total dipole-dipole energy=	0.17972682E+00[Total dipole-dipole energy=	0.14317093E+00[
0.21676078E+01 j	0.0000010000	0.17267229E+01 j	0 147120228.001	
Total charge-dipole+dipole-dipole energy=	0.38395132E+00[0 17745720EL01 1	0.14/13833E+00[
Higher multipole interaction energy =	-0 75325700E+00[Higher multipole interaction energy =	-0 742095728+001	
-0.90847081E+01]	0./3323/001/00[-0.89500967E+01]	0./42000/20100[
Total isotropic repulsion-dispersion=	-0.20073845E+01[Total isotropic repulsion-dispersion=	-0.21696973E+01[
-0.24210199E+02]		-0.26167785E+02]		
Total anisotropic repulsion energy=	0.0000000E+00[Total anisotropic repulsion energy=	0.0000000E+00[
0.0000000E+00]		0.0000000E+00]		
Intermolecular induction energy=	0.0000000E+00[Intermolecular induction energy=	0.00000000E+00[
$\begin{array}{ccc} 0.0000000 \pm 100 \\ D_{1} & D_{2} \\ D_{2} & D_{2} \\ D_{3} & D_{3} \\ D_{$	1001300000000	$D_{\rm V} = 0 0 D_{\rm V} = 0 0 D_{\rm V} = 0 D_{\rm V$	100130000000000	
$PV \in PO(2)$ $P(P) = 0.00E + 00 Pa) \dots \dots$	0.0000000000000000000000000000000000000	rv energy(r = 0.00E+00 Pa)	0.0000000E+00[
Total lattice energy	-0.35783261E+01[Total lattice energy	-0.36992673E+01[
-0.43156649E+02]		-0.44615268E+02]		

II - SOXLEX 03.lem_custom_spli_SOXLEX & 04.lem_custom_nospli_SOXLEX

This example shows how the user makes their own definition of the atomic types, and uses an anisotropic atomatom repulsion potential. This crystal structure was set as one of the challenges in the Fourth Blind Test of Crystal Structure Prediction.² The starting files are trial crystal structures generated by **MOLPAK**, with the dma generated using **GAUSSIAN**³ and **GDMA**.⁴ The repulsion dispersion potential was generated using **SAPT (DFT)**.⁵

A custom potential has different atomic types from FIT or WILLIAMS. The labels file lists all atoms from the input crystal structure file and assigns them the potential types which are used in the dmacrys.mols, dmacrys.dma and potential (pote.dat) files.

The	labels file, defining atomic types
Br2	BR
Br1	BR
Cl	CL
F	FL
C1	CA
C3	CC
C4	CA
C5	CB
H1	НҮ
C6	CD
C2	CB
H2	НҮ

The anisotropic atom-atom repulsion potential requires the definition of a local axis system for each atom, and the anisotropic potential needs to be defined for every pair of atom types:

Starting excerpt from dmacrys.mols file, showing definition of local axis on a Cl atom



The output has

Total isotropic repulsion-dispersion...=-5.8978[-142.2616Total anisotropic repulsion energy....=2.7328[65.9171

which is slightly misleading as when a custom (anisotropic) repulsion is used, all the repulsion, isotropic and anisotropic, comes to 65.9171 kJ/mol and the dispersion contribution is -142.2616 kJ/mol.

NEIGHCRYS automatically sets up the limits used in summing the contributions to the lattice energy. These are a compromise between timing and accuracy and may need changing for different studies. Splines have been introduced rather than abrupt cutoffs in the direct space summations, so to give continuous first derivatives. The effect of splines using the default cutoff (see next example III) is shown, with examples 3 and 4 respectively. The line to bring the splines into play is automatically written out into the *.dmain file.

SPLI 2.0 4.0

This is manually removed from example 4, so that the repulsion-dispersion energy stops abruptly at 15 Å, which is only marginally longer than the c vector. The effect of the splines is expected to be much larger for this molecule than for III or other C/N/O/H molecules, because the dispersion coefficients for Br and Cl are much larger than for first row atoms, their repulsion is longer range, and they have significant atomic quadrupoles, so the inter atomic potentials at the cutoff can be much larger and not necessarily attractive.

II_3 – With splines				II_4 – Without splines
Initial Lattice Energy: Final Lattice Energy:	-84.6923 -85.1785	8 kJ/mol 5 kJ/mol		Initial Lattice Energy: -85.8975 kJ/mol Final Lattice Energy: -86.1616 kJ/mol
a (Ang) beta (Deg) gamma (Deg)	b (Ang)	c (Ang)	alpha (Deg)	a (Ang) b (Ang) c (Ang) alpha (Deg) beta (Deg) gamma (Deg)
Initial => 3.8052 93.7800 90.0000	13.7907	14.5314	90.0000	Initial => 3.8052 13.7907 14.5314 90.0000 93.7800 90.0000
Final => 3.8430 92.5989 90.0000	14.0048	14.5186	90.0000	Final => 3.8286 13.8841 14.6405 90.0000 94.4913 90.0000
F = 8.374435				F = 2.759047
Total run time		••	2.380000	Total run time

III – KAXXAI02 05.largecell_KAXXAI , 06.largecell_KAXXAI_nospli and 07.largecell_KAXXAI_30ang

This example illustrates the use of experimental X-ray structures as input, including when there are two symmetry independent molecules in the asymmetric unit cell. Test 5 uses summation defaults, which are contrasted with other treatments of the direct summation in tests 6 and 7. This molecule, tolfenamic acid, was studied as part of a series of fenamic acids.⁶ The starting file was the crystal structure determination from the CSD, but with the atoms renumbered (to conform to the numbering scheme used in the rest of the series).

Since this is an experimental X-ray structure, the positions of the hydrogen atoms suffer from the systematic foreshortening of bonds to hydrogen. **NEIGHCRYS** has the question

Do you want to standardise bond lengths to hydrogen?

and the response

y(es)

will automatically change the hydrogen atomic positions by elongating the bond lengths to hydrogen to standard neutron values.⁷ This can make a considerable difference to the results of the modelling. The hydrogen atom positions do not need correcting if the crystal structure has been determined by neutron diffraction, or the molecular structure obtained by ab initio optimisation, as the bond lengths to hydrogen should be similar or more accurate.

Although the two molecules in the asymmetric unit are the same, the axis system must be defined for each molecule, and the distributed multipoles calculated separately for each molecule to reflect the differences in conformation. Generating the dma for a two-component system requires generation of each dma separately (using GAUSSIAN³ and GDMA⁴) and the two combined using gdmaneighcrys (supplied with the

DMACRYS/NEIGHCRYS release bundle).

Definition of axis system for two molecules	
MOLX 2 X LINE C F1 2 C F1 5 3	Note that corresponding atoms have been used to define
Y PLANE C F1 2 C F1 5 3 N F1 1 2 X LINE C F1 16 C F1 19 3	the axis system. However, as the molecules differ in the
Y PLANE C_F1_16 C_F1_19 3 N_F1_2 2	central torsion angle and hence centre of mass, there are
	differences in the coordinates of all atoms in the local
	axis system.

Contrasting the . dma entry for two equivalent atoms in the same molecule.								
2 CF12	-0.742683	5.289326	0.071526		32 C F1 16	-0.518412	5.187897	-0.826132
Next 3 Limit 4					Next 33 Limit 4			
-0.059318					-0.063162			
0.002371 -0.008579	-0.044884				0.007418 -0.007997	-0.047917		
-1.263432 -0.020244	-0.019694	0.061451	-0.020968		-1.276015 -0.042648	-0.007954	0.063306	-0.011778
-0.015451 -0.320041	-0.143841	0.004144	-0.033650	-	-0.016989 -0.312273	-0.146901	0.026483	0.015987 -
1.111576 0.161357					1.112546 0.153051			
-1.713140 -0.048270	-0.074806	0.249255	-0.226941		-1.770809 -0.095468	0.051984	0.271955	-0.220172
0.032375 -0.002643					0.021388 -0.011882			
-0.079820 -0.184937					-0.083770 -0.195849			

The molecule itself is quite large, having a distance of 9.79 Å between the two most separated hydrogen atoms. The cutoff over which **DMACRYS** calculates intermolecular interactions is written in the *.dmain file by **NEIGHCRYS**. By default, the limit of the repulsion-dispersion potential is 15 Å, and the higher multipole contributions are calculated for all atoms in molecules whose centre of mass are within 15 Å. The limit of the repulsion-dispersion in the potential file can be greater than this (it is 70 Å in the files provided with the distribution), but it is overridden by the *.dmain. Ideally the summation limits should not be exactly an integer number of lattice vectors.

The three examples given for the lattice energy minimization of KAXXAI02 are (5) with **NEIGHCRYS** defaults of 15 Å cutoffs and splines, (6) with **NEIGHCRYS** default of 15 Å cutoffs, but no splines and (7) with a longer 30 Å cutoff and no splines. Note that since **DMACRYS** does not know that the two molecules in the asymmetric unit cell are the same molecule, all lattice energies need to be divided by 2 to be strictly in kJ/mol and comparable with the lattice energy of the Z'=1 polymorphs of tolfenamic acid.

III_5 – 15 Å cutoff and splines	III_6 – 15 Å cutoff without splines	III_7 – 30 Å cutoff without splines		
CUTO 28.065000 0.534474 RDMA 0.534474	CUTO 28.065000 0.534474 RDMA 0.534474	CUTO 28.065000 1.068948 RDMA 1.068948		
· •				
SPLI 2.0 4.0				
Initial Lattice Energy: -268.4195	Initial Lattice Energy: -267.3524	Initial Lattice Energy: -271.8905		
Final Lattice Energy: -276.3210 kJ/mol	Final Lattice Energy: -275.3161 kJ/mol	Final Lattice Energy: -279.7492 kJ/mol		
a (Ang) b (Ang) c (Ang) alpha (Deg) beta (Deg) gamma (Deg) Initial => 7.6356 11.3050 28.0650 90.0000 93.0300 90.0000 Final => 7.7417 11.5427 27.8747 90.0000 91.2818 90.0000	a (Ang) b (Ang) c (Ang) alpha (Deg) beta (Deg) gamma (Deg) Initial => 7.6356 11.3050 28.0650 90.0000 93.0300 90.0000 Final => 7.7493 11.5629 27.8446 90.0000 91.2531 90.0000	a (Ang) b (Ang) c (Ang) alpha (Deg) beta (Deg) gamma (Deg) Initial => 7.6356 11.3050 28.0650 90.0000 93.0300 90.0000 Final => 7.7436 11.5457 27.8533 90.0000 91.2849 90.0000		
F = 20.042020	F = 22.245230	F = 20.564110		
Total run time 24.500000	Total run time 17.740000	Total run time 104.269997		

When calculating the interatomic interactions over a larger number of molecules (example 7 with 30 Å limits to interaction calculation), the accuracy of the calculation is much greater as the longer range interactions are included, but the time taken is much longer. Increasing the cutoff or adding splines usually makes the lattice energy slightly lower because the attractive dispersion is the longest range contribution.

IV - KONTIQ01 08.2comp_KONTIQ

This illustrates using two different molecules in the asymmetric unit, with one being water. Gallic acid monohydrate was set as one of the challenges in the Fifth Blind Test of Crystal Structure Prediction,⁸ to predict the structures of two new polymorphs, with this structure and that of another polymorph already published. It also had the complication of being a hydrate. Further screening resulted in characterising a further monohydrate, i.e. five polymorphs in all.⁹

The starting files in this example are from the end point of a **CrystalOptimizer** calculation, which optimizes the molecular conformation and crystal lattice together, including generating the .dma using **GAUSSIAN**³ and **GDMA**⁴ for sufficient conformations that a database can be used to speed these calculations over a large set of crystal structures.¹⁰ Note that **gdmaneighcrys** has to be used as part of the process to combine the gallic acid and water dmas.

DMACRYS automatically assigns hydrogen atoms in water molecules as a different atomic type – H_Wa. The *exp-6* potential used for the water hydrogens was the same as for other polar hydrogen atoms in this case, but a different

potential could have been chosen for H_Wa without the need to use a custom potential and labels file. Since the pote.dat *exp-6* input is required for each pair of atomic types, the **crosspot** utility supplied with the **DMACRYS/NEIGHCRYS** release bundle is useful for generating the unlike interaction parameters using the combining rules for extensions to the *exp-6* potentials of the FIT or Williams type.

V – FINVAZ 09.salt_FINVAZ & 13.properties_FINVAZ

Amantidine hydrochloride is an example of a lattice energy minimization, using the Williams potential, of a salt containing a spherical ion. In this case, the dma is generated for the molecular ion using GAUSSIAN³ and GDMA⁴ and the resulting dmacrys.dma file is manually edited to add the unit negative charge for the chloride ion. Rotations of the spherical chloride ion are meaningless; hence it does not require an axis system etc. It is also an example of a centred cell, C2/c (Z=8, Z'=1).

Salts such as FINVAZ have lattice energies that are orders of magnitude more stabilising than neutral systems, because of the high electrostatic interaction between the two species in the crystal. Nonetheless, as the properties calculation shows, such structures still have phonon modes of comparable frequencies to neutral organic molecules.

VI – Ephedrinium Tartrate 10.largesalt_WEMGEK

Ephedrinium Tartrate is provided as an example of a lattice energy minimization of a salt with two large molecular components. It is a structure where the hydrogen atoms were positioned following a CSP study.¹¹ The same considerations as in previous examples should be made, namely it has large lattice energy, and the dma needs to be calculated separately for the two species and combined using gdmaneighcrys.

VII – CBMZPN10 11.properties_CBMZPN

DMACRYS can calculate the second derivative properties of the crystal at the lattice energy minimisation, namely the elastic constants¹² and k=0 phonon frequencies.¹³ This is done with accurate second derivatives of the lattice energy, not the estimates made by updating the Hessian matrix used within the lattice energy minimisation. (These estimated second derivatives are used to calculate an approximate elastic constant matrix and eigenvalues of the Hessian to test whether the Born stability criterion is met, or whether the structure is a transition state, see **16.symmred_PAPTUX**). The elastic tensor is usually aligned so that *z* is along *c*, *x* is parallel to *a* and *y* is in the *ab* plane, but please check fort.21 for higher symmetry space groups. The lowest eigenvalue and eigenvector of the shear submatrix is calculated as a guide to whether the crystal has a particularly weak plane. A range of properties of a microcrystalline aggregate, by various approximate averaging procedures,¹⁴ are also reported.

The example of a properties calculation with the FIT potential is for the most stable polymorph of carbamazepine (form III). The example files here are a lattice energy minimum found in a rigid molecule CSP study of carbamazepine.

In order to calculate second derivative properties, it is usual to start at a lattice energy minimum (fort.16 from a previous standard lattice energy minimisation run) to save the time in minimising with the more expensive and accurate 2^{nd} derivatives calculation. This is why in the example output the lattice energy changes so very little in the fort.12 and why the goodness of fit (F) is so low.

To run a properties calculation, following a **NEIGHCRYS** run starting from an estimated lattice energy minimum, the *.dmain must be edited to change the line that said "STAR PLUT" to say "STAR PROP" and remove the NOPR directive so that the program uses the most accurate second derivatives. It is also necessary to increase the accuracy of the Ewald summation by adding the line ACCM 100000000 in the section of the *.dmain file that deals with changes to default parameters.

Output from pr	operties calculations	
Zone Centre Phon THz	on Frequencies cm-1	The first 3 frequencies should be zero, to within
0.0000 0.0000 0.0000 0.0000	0.0000 0.0000 0.0000	numerical error.

1.2704 42.3756 1.3652 45.5384 1.4234 47.4784
1.5027 50.1256 1.6125 53.7867
1.6128 53.7961 1.7326 57.7942
2.0264 67.5939 2.0904 69.7281
2.2966 76.6065 2.4722 82.4652
2.5321 84.4618 2.6035 86.8428
2.7487 91.6871 3.0024 100.1495
3.0663 102.2815 3.1597 105.3979
3.2436 108.1949 3.8337 127.8800
3.9112 130.4624 4.0285 134.3770
MECHANICAL PROPERTIES CALCULATIED FROM HESSIAN MATRIX
Elastic stiffness tensor (Cij): GPa (10**10 DYNE/CM**2)
$\begin{array}{cccccccccccccccccccccccccccccccccccc$
0.00000 0.00000 0.00000 5.09660 0.00000 $-0.637033.77841$ 0.32893 -0.69397 0.00000 11.35118 0.00000
-3.47641 0.32933 -0.00397 0.00000 11.3110 0.00000 11.39759
Born stability criteria satisfied.
Corresponding eigenvector 0.99330 0.00000 0.11557)
Elastic compliance tensor (Sii): GPa^(-1) (10**-10 CM**2/DYNE)
0.11605 - 0.02046 - 0.07303 0.00000 0.03176 0.00000 - 0.02046 0.03528 0.00480 0.00000 - 0.00700 0.00000
-0.07303 0.00480 0.12111 0.00000 -0.01522 0.00000 0.00000 0.00000 0.00000 0.16776 0.00000 0.00938
0.03176 -0.00700 -0.01522 0.00000 0.09711 0.00000 0.00000 0.00000 0.00000 0.00938 0.00000 0.08826
Mechanical properties calculated for a microcrystalline aggregate.
Bulk modulus from elastic constants (Voigt averaging): 11.61348 GPa Bulk modulus from compliance constants (Reuss averaging): 10.52026 GPa
Arithmetic mean: 11.06687 GPa Geometric mean: 11.05336 GPa
Shear modulus from elastic constants (Voigt averaging): 8.57280 GPa
Shear modulus from compliance constants (Reuss averaging): 5.99050 GPa Arithmetic mean: 7.28165 GPa
Geometric mean: /.1662/ GPa
Young's modulus from elastic constants (Volgt averaging): 20.63980 GPa Young's modulus from compliance constants (Reuss averaging): 15.10453 GPa
Geometric mean: 17.65657 GPa
Poisson's ratio from elastic constants (Voigt averaging): 0.20380 GPa
Arithmetic mean: 0.23225 GPa Geometric mean: 0.23050 GPa
Single crystal Young's modulus along optical axes x-component: 8.61669 GPa
y-component: 28.34411 GPa z-component: 8.25685 GPa
arithmetic mean: 15.07255 GPa

Experimental elastic constants and phonon frequencies are often not available, and the main reason for calculating these properties is to estimate the zero-point vibrational energy and thermal contribution to the Helmholtz free energy. In this example, the resulting *.dmaout file has been automatically analysed to extract the phonon frequencies¹³ and elastic constants.¹² More recent work by Nyman and Day^{15, 16} has developed the use of **DMACRYS** to converge free energy differences between polymorphs, by using supercells to sample the Brilloiun zone. Their utilities **AutoLD**, and **AutoFree** for doing such calculations are included in the release bundle.

<pre>kD(x) bit = -1.91265860137880 entropy from acoustical modes (J/molK) : from Debye fiunction bit (J/molK) = -2.07364893717513 from other bit (J/mOlK) = -15.7096067768345 Total Entropy (J/molK) = -104.997925394349</pre>						
thermal energy worked out for T (K) = 298.00000000000						
thermal E from optical modes (kJ/mol) = -15.3634642623644 thermal E from -kTD(x) (kJ/mol) = -0.569972263210882 thermal E from second bit (kJ/mol) = -3.01950438657552 Total Thermal Energy (kJ/mol) = -18.9529409121508						
Vibrational Free Energy (including ZPE): Fvib (kJ/mol) = -16.2295901832753						

VIII – DCLBEN06 12.properties_DCLBEN

This is an example of a properties calculation for the β triclinic P-1 Z'=0.5 polymorph of dichlorobenzene with a custom potential using a non-empirically derived atom-atom repulsion-dispersion potential for anisotropic chlorine.¹⁷ This illustrates the more accurate type of potential that is likely to be needed for studying properties which are very sensitive to the 2nd derivatives, as shown by the comparison with the recorded low temperature spectrum.¹⁸

VIII_12 from KC	ONTIQ_A23	12.dmaou	ut file	Experimental room temperature spectrum ¹⁸		
Zone Centre Phon THz	on Frequencies cm-1				cm ⁻¹	
0.0000	0.0000				Three zero frequency modes, translating entire crystal	
0.0000 0.0000	0.0000				56	
1.4722	49.1063				65	
3.1102	103.7450				103 – Fortuitously close!	
Only one moled	cule in the un	it cell gives	s few lattice m	nodes,	but the triclinic cell has a full elastic tensor	
Elastic stiffnes	s tensor (Cij)	: GPa (10**	10 DYNE/CM**2)	0.047		
23.58876 9.62432	9.62432 17.06762	6.02847	2.05584	-2.341	157 -3.50538 018 -3.71365	
6.02847	6.37537	13.46004	-1.47659	0.102	214 0.10800	
-2.34157	∠.US584 -0.82018	-1.4/659	4.2//30	-0.558	520 -2.35944 188 2.78290	
-3.50538	-3.71365	0.10800	-2.35944	2.782	90 9.11696	

IX - 4-fluorotoluene 14.press_YICDIZ

The starting file in this example is a crystal structure determination where the liquid was forced to crystallize at high pressure and the X-ray data collected under those extreme conditions.¹⁹

Following the NEIGHCRYS run, the YICDIZ.dmain needs to be edited to add a line to include pressure. "PRES 0.5 GPa" after the CONP directive.

The resulting lattice energy from the fort.12 is very high because it includes the PV term.

Initial Lattice Energy: -11.4386 kJ/mol Final Lattice Energy: -11.9699 kJ/mol

X – OBEQUIJ 15.induction_OBEQUJ

This illustrates the calculation of induction energy (sometimes called polarization energy, as we are using the longrange part of the induction energy, as calculated from point polarizabilities and multipoles) and lattice energy minimisation, including this term. The inclusion of damped induction and dispersion up to C₁₀ are new features of **DMACRYS** for research work using non-empirical potentials.

This molecule was set as one of the challenges in the Fifth Blind Test of Crystal Structure Prediction.²⁰ The example uses the ab initio optimised molecular structure and has a non-empirical potential (c.f. SOXLEX⁵ derived using CamCASP²¹) which has been fitted excluding the induction energy. The modelling of the induction energy using a distributed dipolar polarizability model¹ is the novel feature of this example. The molecular structure, atomic multipole moments, dipolar polarizability tensors and atom-atom isotropic C₆ dispersion coefficients were derived from the PBEO exchange-correlation functional and the Sadlej pVTZ basis set with the Tozer-Handy asymptotic correction. The CUSTOM repulsion-dispersion model assumed that the parameters were transferable between all

hydrogens and the four carbon atoms bonded to hydrogen (see labels file). This proved to be a lousy model potential, giving rise to considerable slippage in the *ab* plane, from deficiencies in the repulsion-dispersion potential. The *c* vector is large, so a larger cutoff should be used, but this significantly increases the computer time.

Note that the induction energy has to be iterated to convergence, so there is numerical noise in taking numerical derivatives. Hence the induction energy is only evaluated at a given structure. (i.e. in an optimisation, the induction energy is evaluated at the initial and final structure, and the forces from the induction energy are not used in the minimisation.)

The distributed dipolar polarizabilities are in the file dmacrys.dma.pol, which starts:

```
# Static polarizabilities
   C_C1_1_
2.46610
                       -0.277986 -1.883193 -0.000535
                                                            RANK 1
   0.00000 14.52040
                       12.08170
   0.00000
               4.63450
        C2 2
                        0.816974 0.674908 -0.000979
                                                            RANK 1
 2
    2.67470
           23.67240
    0.00000
               4.44500
                        17.25500
```

This is automatically read-in by **NEIGHCRYS** provided that the need for induction is specified near the start of the dmacrys.mols file by specifying the induction damping coefficient DIND, for example

```
      MOLX 1

      X LINE C_C2_2
      N_N1_1
      1

      Y PLANE C_C2_2
      N_N1_1
      1 0_01_1
      2

      DIND 1.0885
      N_NS
      C_C1_1
      2

      Z LINE C_C1_1
      0_01_1
      1
      1 C_C2_2
      1

      X PLANE C_C1_1
      0_01_1
      1
      1 C_C2_2
      1
```

In the *.dmaout file, the polarizabilities are written out, and the phrase "Induction damping has been requested" precedes the damping parameters. The induction calculation at the starting structure reports

First order indu Atom Label C_C1_1 C_C1I1417 C_C2_22	ced moments Q10 -0.003618019 0.003659108 -0.008601462	(atomic units) Qlc 0.011915082 0.012115435 0.138004038	Q1s 0.075723632 0.075643569 0.154284011	
$\begin{array}{c} N & N1 & 1 & 7 \\ N & N1 & \overline{114} & 7 \\ N & N2 & 2 & 8 \\ N & N2 & 15 & 66 \\ O & 01 & 1 & 9 \\ O & 01 & 14 & 73 \\ \end{array}$	-0.009497820 0.009513033 -0.024434997 0.024418822 0.005452855 -0.005383297	0.016845367 0.016922579 -0.047180778 -0.047173454 -0.012324706 -0.012226963	0.049625221 0.049644818 0.040711153 0.040708845 0.096145953 0.096193965	
H H1 4 13 H H1 117 101 FIrst-order indu Converged induce Atom Label C C1 1 1 C C1 1 1 C C2 2 2	-0.003580690 0.003580690 ction energy d moments (at Q10 -0.006422005 0.006461885 -0.013724200	-0.001161334 -0.001161334 -6.72028888 tomic units) Qlc 0.032483996 0.032705296 0.185647096	-0.005492413 -0.005492413 kJ/mol. Q1s 0.128535480 0.128461236 0.218022713	
. N1 1 7 N_N1114_59_ N_N2 2 8 N_N2115 66 O_O1 1 9 O_O1114_73_	-0.011628952 0.011638417 -0.023572019 0.023529057 0.010599723 -0.010532090	0.039485056 0.039596528 -0.050642876 -0.050630339 -0.011842498 -0.011733757	0.071653984 0.071663080 0.062891544 0.062890475 0.151069302 0.151117136	
H H1 4 13 H H1 17 101 Induction energy	-0.002437387 0.002440094 converged to	0.001578367 0.001584326 -9.64709323	-0.004484172 -0.004475367 3 kJ/mol with	17 iterations.

The induced dipole moments should be modest, reflecting the atoms which are most polarizable, with the induction energy being stabilising and converge in a reasonable number of iterations. The induction energy is written out as a component of the initial and final lattice energy. The initial energies and induced moments should be well reproduced, but minor discrepancies in the final structure and energy may just reflect numerical noise.

XI – PAPTUX 16.symmred_PAPTUX

Lattice energy minimisation within the constraints of space group symmetry may lead to a structure which is not a true minimum, as detected by considering the eigenvalues of the second derivative matrix (~Born stability criterion). When this happens, then another run in a lower symmetry space-group, usually with more independent molecules in the asymmetric unit cell is required. This example shows the ability of **NEIGHCRYS** to generate the required new input files and run in lower symmetry.

The anti-inflammatory naproxen was studied because the structure of the marketed enantiopure form was known but not that of the racemate.²² The crystal energy landscape showed that the racemic form was more stable, but experimentally it proved impossible to obtain single crystals, so this structure had to be determined from powder Xray diffraction. A good match to the experimental data was obtained by the global minimum of the lattice energy landscape. However this Z'=1 structure had a negative eigenvalue of the Hessian, showing that it was a saddlepoint. This example demonstrates how the symmetry was lowered to give a true minimum, requiring a second run of **NEIGHCRYS** and **DMACRYS** to give a Z'=2 structure with the small energy lowering (~ 1 kJ/mol). Solid state NMR was used to confirm that structure was indeed *Pbca* Z'=1, confirming the estimate that even the zero-point motions would average over the Z'=2 *Pca2*₁ lattice energy minima.

The files in the directory reason are the original **NEIGHCRYS** and **DMACRYS** runs, showing the fort.12 detects that the structure is not a true minimum.

Warning - Non-zero eigenvalues found. A subgroup has a lower energy structure

To find the correct eigenvalue (~symmetry element) to remove, the PAPTUX_CO1.dmaout must be inspected. The relevant section looks like this.

EIGENVALUES FOR	REPRESENTATION	1		
0.54740738E+01	0.10484391E+02	0.21314702E+02	0.28925942E+02	0.32695686E+02
0 58755977E+02				
ETCENNATUES FOR	DEDDECENTATION	2		
LIGENVALUES FOR	REPRESENTATION	2 000000017.01	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0 040107567.00
0.46788755E-13	0.41550997E+01	0.88608361E+01	0.26927779E+02	0.34219756E+02
0.40828089E+02				
EIGENVALUES FOR	REPRESENTATION	3		
-0.13293712E+01	0.20725220E-13	0.84737284E+01	0.15625808E+02	0.22016299E+02
0 55803587E+02				
FTCENUATIES FOR	DEDDEGENITATION	1		
CIGENVALUES FOR	A DCAADEOOD OO		0.0000000000000000000000000000000000000	0 (204401(7)00
0.80397293E+01	U.26443582E+U2	0.30901556E+02	0.39325824E+U2	0.63044816E+02
0.74113852E+02	0.99367508E+02			
EIGENVALUES FOR	REPRESENTATION	5		
-0.91065363E-14	0.10675599E+02	0.14989850E+02	0.24870560E+02	0.29247885E+02
0.93156502E+02				
EIGENVALUES FOR	REPRESENTATION	6		
0.016400620101	0.244049555.02	0 201017270.02	0 20627/11100	0 502000720002
0.91040002E+01	0.2440465555+02	0.39191/2/6+02	0.3903/4116+02	0.303800/26402
0./9/33633E+02	0.920816/9E+02	_		
EIGENVALUES FOR	REPRESENTATION	7		
0.75603086E+01	0.20745637E+02	0.32486537E+02	0.42178187E+02	0.48736699E+02
0.73609246E+02	0.76396408E+02			
EIGENVALUES FOR	REPRESENTATION	8		
0 93971337F+01	0 20082801F+02	0 28845175E+02	0 394086935+02	0 45319387E+02
0.711404600.00	0.114016410.00	0.100000000000	0.307000935102	0.30010102
0./11404696+02	0.11491641E+03	0.15938635E+03	0.38/29/56E+03	

In the next **NEIGHCRYS** run, the eigenvalue that is negative should be removed. In this case, the eigenvalue for representation *3* is negative and is the relevant one. (There are always 3 eigenvalues that are zero to within numerical error, for translation of the entire crystal. These are not exactly zero and so may also be negative, and are here in representations 2, 3 and 5). It can sometimes be the case that there are 2 negative eigenvalues, and it is recommended that you remove the larger magnitude one first, reminimize the structure, check the output, and run **NEIGHCRYS** again to remove the other if necessary. You should never need to remove the last representation on the list as this is the totally symmetric representation.

The subsequent run (**16.symmred_PAPTUX**) of **NEIGHCRYS** is distinguished by using this value 3 to answer the questions

Input zero for no symmetry subgroup or n to remove representation n 3 CVECTOR 5.51670 Symmetry reduction for representation 3 complete input zero to end symmetry reduction or n to remove new representation n 0 which produces the files with parameters for two independent naproxen molecules: dmacrys.new.mols as a replacement to be used instead of the dmacrys.mols and PAPTUX_CO1.new.dma instead of dmacrys.dma. These files are needed as input if it proves necessary to remove another representation. However, in this example, at the end of the run, a true minimum has been reached, with the fort.12 summary file stating:

Valid minimisation - converged on gradients Eigenvalue calculation is exact Minimum passed eigenvalue test

Change in Hessian matrix during run 16.symmred_PAPTUX

At the start of the run with reduced symmetry, we have the negative eigenvalue appearing in the totally symmetric representation

EIGENVALUES FOR RE	PRESENTATION	1			
0.54749828E+01	0.91665656E+01	0.10486225E+02	0.21324804E+02	0.24408438E+02	
0.28940038E+02	0.32700849E+02	0.39198516E+02	0.39656679E+02	0.50404872E+02	
0.58785206E+02	0.79773166E+02	0.92096847E+02			
EIGENVALUES FOR R	REPRESENTATION	2			
-0.19184654E-12	0.41556496E+01	0.80411707E+01	0.88624634E+01	0.26448294E+02	
0.26940703E+02	0.30916946E+02	0.34226134E+02	0.39332327E+02	0.40848079E+02	
0.63075419E+02	0.74150025E+02	0.99385942E+02			
EIGENVALUES FOR R	REPRESENTATION	3			
0.28421709E-13	0.75617577E+01	0.10677549E+02	0.14992172E+02	0.20749220E+02	
0.24874684E+02	0.29262150E+02	0.32502922E+02	0.42184663E+02	0.48760752E+02	
0.73621962E+02	0.76433558E+02	0.93202329E+02			
EIGENVALUES FOR R	REPRESENTATION	4			
-0.13299736E+01 -	-0.14210855E-13	0.84752917E+01	0.93988285E+01	0.15629116E+02	
0.20086029E+02	0.22019776E+02	0.28859255E+02	0.39428316E+02	0.45327300E+02	
0.55830304E+02	0.71153052E+02	0.11497258E+03	0.15941309E+03	0.38736229E+03	
1. 1. C					
10 700 7000 1000	\mathbf{t} of the element	n n n n	atry radiiction	wa hava a trua	minimum (only 2 zoro oigonyaluoc)
In the final repor	t of the eigenv	values in symme	etry reduction,	we have a true	minimum (only 3 zero eigenvalues)
EIGENVALUES FOR RE	t of the eigenv	values in symme	etry reduction,	we have a true	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01	T OF THE EIGENN PRESENTATION 0.90017777E+01	/alues in symme 1 0.10909412E+02	0.21923355E+02	0.24717232E+02	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02	T Of the eigenv PRESENTATION 0.90017777E+01 0.29355767E+02	/alues in symme 1 0.10909412E+02 0.37714874E+02	0.21923355E+02 0.39091498E+02	0.24717232E+02 0.51114580E+02	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02	T Of the eigenv PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02	0.21923355E+02 0.39091498E+02	0.24717232E+02 0.51114580E+02	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R	T Of the eigenv EPRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 EPRESENTATION	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2	0.21923355E+02 0.39091498E+02	0.24717232E+02 0.51114580E+02	minimum (only 3 zero eigenvalues)
In the final report EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.600161.01E+02 EIGENVALUES FOR R 0.17763568E-13	C Of the eigenv EPRESENTATION 0.29017777E+01 0.29355767E+02 0.78624602E+02 VEPRESENTATION 0.43762867E+01	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01	0.21923355E+02 0.39091498E+02 0.90276714E+01	0.24717232E+02 0.51114580E+02 0.27077689E+02	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02	C Of the eigenv PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 WEPRESENTATION 0.43762867E+01 0.30606861E+02	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02	we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02	CONTRACTOR CONTRAC	/alues in symme 1 0.10909412E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02 0.10569275E+03	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02	we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02 EIGENVALUES FOR R	TOT THE EIGENV PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 0.43762867E+01 0.30606861E+02 0.74740580E+02 0.74740580E+02 0.74740580E+02	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02 0.10569275E+03 3	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02	we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02 EIGENVALUES FOR R 0.42632564E-13	T Of the eigenv PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 KEPRESENTATION 0.43762867E+01 0.30606861E+02 0.74740580E+02 0.74740580E+02 0.80571557E+01	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02 0.10569275E+03 3 0.10755696E+02	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02 0.15344159E+02	<pre>we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02 0.21064826E+02</pre>	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02 EIGENVALUES FOR R 0.42632564E-13 0.23683924E+02	T Of the eigenv PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 KEPRESENTATION 0.43762867E+01 0.30606861E+02 0.74740580E+02 KEPRESENTATION 0.80571557E+01 0.28396204E+02	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02 0.1055696E+02 0.34279322E+02	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02 0.15344159E+02 0.38682797E+02	<pre>we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02 0.21064826E+02 0.49266106E+02</pre>	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02 EIGENVALUES FOR R 0.42632564E-13 0.23683924E+02 0.69790339E+02	TOT THE EIGENY PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 WEPRESENTATION 0.43762867E+01 0.30606861E+02 WEPRESENTATION 0.80571557E+01 0.28396204E+02 0.73566111E+02	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02 0.105569275E+03 3 0.10755696E+02 0.34279322E+02 0.92188150E+02	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02 0.15344159E+02 0.38682797E+02	<pre>we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02 0.21064826E+02 0.49266106E+02</pre>	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02 EIGENVALUES FOR R 0.42632564E-13 0.23683924E+02 0.69790339E+02 EIGENVALUES FOR R	T Of the eigenv PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 @EPRESENTATION 0.43762867E+01 0.30606861E+02 0.74740580E+02 @EPRESENTATION 0.80571557E+01 0.28396204E+02 @CRESENTATION 0.73566111E+02 @EPRESENTATION	/alues in symme 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02 0.10569275E+03 3 0.10755696E+02 0.34279322E+02 0.92188150E+02 4	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02 0.15344159E+02 0.38682797E+02	<pre>we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02 0.21064826E+02 0.49266106E+02</pre>	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02 EIGENVALUES FOR R 0.42632564E-13 0.23683924E+02 0.69790339E+02 EIGENVALUES FOR R -0.56843419E-13	T Of the eigenv PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 WEPRESENTATION 0.43762867E+01 0.30606861E+02 0.74740580E+02 WEPRESENTATION 0.80571557E+01 0.28396204E+02 0.73566111E+02 WEPRESENTATION 0.28814824E+01	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.3087065E+02 0.10755696E+02 0.34279322E+02 0.92188150E+02 4 0.88057206E+01	0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02 0.15344159E+02 0.38682797E+02 0.91321059E+01	<pre>we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02 0.21064826E+02 0.49266106E+02 0.16240193E+02</pre>	minimum (only 3 zero eigenvalues)
In the final repor EIGENVALUES FOR RE 0.62821481E+01 0.25674871E+02 0.60016101E+02 EIGENVALUES FOR R 0.17763568E-13 0.27322200E+02 0.58616835E+02 EIGENVALUES FOR R 0.42632564E-13 0.23683924E+02 0.69790339E+02 EIGENVALUES FOR R -0.56843419E-13 0.19083826E+02	TOT THE EIGENY PRESENTATION 0.90017777E+01 0.29355767E+02 0.78624602E+02 WEPRESENTATION 0.43762867E+01 0.30606861E+02 0.74740580E+02 WEPRESENTATION 0.80571557E+01 0.28396204E+02 0.73566111E+02 WEPRESENTATION 0.28814824E+01 0.20499881E+02	/alues in symme 1 0.10909412E+02 0.37714874E+02 0.95300860E+02 2 0.88027414E+01 0.33087065E+02 0.10755696E+02 0.34279322E+02 0.92188150E+02 4 0.88057206E+01 0.28225705E+02	etry reduction, 0.21923355E+02 0.39091498E+02 0.90276714E+01 0.35887132E+02 0.15344159E+02 0.38682797E+02 0.91321059E+01 0.39943200E+02	<pre>we have a true 0.24717232E+02 0.51114580E+02 0.27077689E+02 0.42707633E+02 0.21064826E+02 0.49266106E+02 0.16240193E+02 0.16240193E+02 0.42124144E+02</pre>	minimum (only 3 zero eigenvalues)

Note that in symmetry reduction, although the resulting structure is Z'=2, the lattice energy has not been doubled, but is only slightly lower (i.e. lattice energy at the saddle point is -134.1608 kJ/mol (Final energy in initial minimisation in reason) and at the minimum is -134.7353 kJ/mol). This is done so that automated symmetry reduction, as part of a CSP study, keeps all the structures on the same energy scale.

XII – TEVSOD 17.pasting_TEVSOD

The **NEIGHCRYS** / **DMACRYS** combination is used for minimizing experimental crystal structures holding the conformation of the molecule rigid throughout. It is often desired to minimize a crystal structure with an alternative conformation of the molecule, such as the ab initio optimized conformation that is being used as input into a Crystal Structure Prediction study. **NEIGHCRYS** was written with the ability to change to a slightly different molecular structure within a crystal structure by aligning the molecule fixed axes of the new and original molecule. This will not work if there is a significant change in the molecular conformation, or the molecules are not identical. The use of the separate utility **optimalpaste** to produce a crystal structure which optimises the overlay of the two molecules is recommended in preference to this option.

The example chosen for this is Cyheptamide. Before beginning, **NEIGHCRYS** is run to set up the molecule in the local axis system. Text from the fort.21 is then used for the **GAUSSIAN** input file, and **GAUSSIAN** is run specifying that the output should be written in atomic units. Alternatively, the output from **GAUSSIAN** in Angstroms can be used, but converted to Bohr with a utility such as http://bast.fr/angstrom-bohr/ The coordinates of the optimized geometry in atomic units are put into the coord_to_paste_au file with the atom labels generated by **NEIGHCRYS**. EVERY molecule in the crystal structure needs to appear in this file (with the sign of the Z coordinates inverted in inverted molecules), so it will look very similar to the end of the fort.21 file.

When **NEIGHCRYS** is run pasting in a new molecule, the end of the fort.21 file is first written out with the original coordinates and then with the new pasted coordinates.

I – AXOSOW 18.defaults_AXOSOW

This example uses the same **NEIGHCRYS** input files as for **02.lem_will01_AXOSOW** but all the filenames have been changed. When **NEIGHCRYS** is run, d should be the answer to the first question, and most of the rest of the questions do not need any answers as the filenames are picked up automatically.

XIII – BT_XXV **19.largecubic_BTXXV**

This is example is of a large cell, where Z is 24. The cubic space group of Pa-3 caused problems with previous version of **NEIGHCRYS**, particularly as the two-component system is so large.

XIII – BT_XXV 20.largecubic_HESS_BTXXV

This is an example of reuse of the hessian calculation from a previous run. Both example 19 and example 20 use exactly the same *.dmain, except that it has been edited to include the "HESS 1" directive in example 20, which instructs **DMACRYS** to write out the hessian at the end of the minimization, and use this file, if it exists, at the starting point of a new minimization. In this example, the dmahessian file is present in the files for example 20, which enables **DMACRYS** to set up the minimization using this hessian and so it runs with fewer iterations, and achieves exactly the same result for the minimization.

Results from 19.largecubic_BTXXV (no us	e of hessian)	Results from 20.largecubic_HESS_BTXXV (reuse of		
		hessian)		
Number of Iterations: 25 Lattice is P Centred. Calculated Lattice En Initial Lattice Energy: -206.0935 kJ/mo Final Lattice Energy: -206.4054 kJ/mo F = 0.494052	ergies 1 1	Number of Iterations: 10 Lattice is P Centred. Calculated Lattice Energies Initial Lattice Energy: -206.0935 kJ/mol Final Lattice Energy: -206.4054 kJ/mol F = 0.494052		
Timing information		Timing information		
Program segment seconds)	Time taken (CPU	Program segment seconds)	Time taken (CPU	
Time to set things up Reciprocal space part of Ewald sum Real space part of Ewald sum Short range potential calculation	4.090000 29.169983 1.940063 0.419922	Time to set things up Reciprocal space part of Ewald sum Real space part of Ewald sum Short range potential calculation	4.000000 29.099991 1.930038 0.399963	
Higher multipole energy calculation		Higher multipole energy calculation		
Energy calculation First derivative chain rule Second derivative chain rule All other program sections Total run time	76.470276 1.339844 49.269653 472.390285 635.090027	Energy calculation First derivative chain rule Second derivative chain rule All other program sections Total run time	76.739655 1.380096 49.389481 130.330765 293.269989	

This shows that the energies are exactly the same at the start and end of the minimization. There is in fact no difference in the output fort.16 files. However, the time saving is made in reducing the number of steps required to minimize the structure, and the overall saving is over 50%.

Normally, the HESS directive is useful when the user is starting an optimization with a slightly different *.dmain file, e.g. with a slightly different molecular conformation, a slightly different dma, or a different cutoff, or to restart an optimization.

XIV – BT_XXIII 21.longcell_BTXXIII

In this example, an unconventional cell setting has been used with an acute angle, resulting in the molecule spanning more than 3 unit cells. **NEIGHCRYS** does not assign the atoms to the same molecule if this happens. This is seen in the "reasons" folder, in the section of the fort.21 headed "Inequivalent basis atoms", where O_01_2 is not included in molecule 1.

at ind	om ex 1 2	atomic number 6 6	na C_F1_1 C_F1_2	12	input name C1 C2	molecule number 1 1	invert flag F F
I	39 40 41 42 43	7 8 8 17 17	N_F1_1_ O_F1_1_ O_01_2_ CIF1_1_ C1F1_2_	39 41 40 42 43	N1 02 01 C11 C12	1 1 0 1 1	म म म

NEIGHCRYS is run in "extended mode" in this example, answering all the additional questions with the default values except

What value of MAXSCH do you want to use to locate whole molecules? [3] $_{4}$

The fort.21 now has all the atoms in the molecule, and **DMACRYS** can run.

Note that this is not a real crystal structure, but a search-generated structure for the Blind Test, and so has many problems in the crystal structure itself. It is recommended crystal structures are reset to the conventional unit cell using a program such as **PLATON** (<u>http://www.cryst.chem.uu.nl/spek/platon/</u>) or **ISOCIF** (<u>http://stokes.byu.edu/iso/isocif.php</u>).

XV – BT_XXII 22. damped_C6_dispersion_BTXXII

In this example, the damping parameter is included in the dmacrys.mols file. Note that the repulsion-dispersion potential is of the DBUC type, and so has $C_8^{\prime\kappa}$ and $C_{10}^{\prime\kappa}$ terms for each interaction, which are zero.

The dmacrys.mols file looks like:

MOLX 1 X LINE C A1 1 C A6 4 3 Y PLANE C $\overline{A1}$ 1 C A6 4 3 C A2 2 1 DDMP 2.47349488

and the start of the pote same as CP.dat looks like:

 DBUC
 C_A1
 C_A1

 3.83545
 0.3533404913
 0.1704323228
 0.0
 0.0
 0.0
 70

 ENDS
 DBUC
 C_A1
 C_A2
 3.83431
 0.3752002353
 0.2174671841
 0.0
 0.0
 0.0
 70

 ENDS
 0.3752002353
 0.2174671841
 0.0
 0.0
 0.0
 70

III – KAXXAI 23.largecell_KAXXAI_30ang

This is exactly the same minimization as example 07, but rather than manually editing the *.dmain file, the larger cutoff is set with the extended question set of **NEIGHCRYS**. The extra question that is not answered with the default answer is:

What is the desired cutoff in Angstroms [15.0]

XVI – PYRDIN 24.dampedispersion_PYRDIN and 25.polarizability_PYRDIN

These illustrate the use of a non-empirical potential for pyridine in the solid state²³ where the model potential had been parameterised by using a state-of-the-art methods for defining the atomic multipoles, polarizabilities and dispersion coefficients and fitting the short range terms adequately required addition forms of anisotropic repulsion.²⁴

As in example 22, the damped dispersion parameter is included in the dmacrys.mols file, and DBUC type repulsion-dispersion potentials are used, and anisotropy is included, for example:

0 0 2	0 2 0 0	0 0 2 2	2 2 0 0	2 0.066 2 -0.086 2 0.066 2 -0.086	35000 556300 35000 556300				
ENDS	NT NT	1	0 01						
DROC	N N	L	0_01						
650	.973	954	0.3	30405000	14.6176	92.5942	1046.8489	0.00	70.00
ANIS	N N1	L	C C1	_					
0	0	0	1	1 0.069	24000				
0	1	0	1	1 0.017	77100				
0	0	1	0	1 0.055	518400				
1	0	1	Ó	1 -0.009	61700				
Ō	õ	Ô	2	2 -0.040	85600				
õ	2	õ	2	2 -0 060	18000				
0	0	2	0	2 0.000	25000				
0	0	4	0	2 0.000	55000				
	0	2	0	2 -0.086	56300				
ENDS									

Furthermore, example 25 includes the induction energy, which is initiated through use of the DIND 1.25 directive giving the induction damping coefficient in the dmacrys.mols file and the required dmacrys.mols.pol file with the polarizabilities.

Comparing the output from the optimisations **24.dampedispersion_PYRDIN** and **25.polarizability_PYRDIN**, it is seen that the unit cell is no different, but that the energy is different. This is because it is not possible to carry out the perfect lattice calculation in **DMACRYS** including the forces due to induction. (Expert use has shown that this approximation has very little effect on the structure.²⁵) The lattice energy minimization part of the calculation is carried out neglecting the forces from the induction. However, **DMACRYS** calculates the induction energy both at the start and at the end, and adds this term to the total initial and final lattice energy.

<pre>25.dampedispersion_PYRDIN PYR1opt.dmaout first</pre>	25.polarizability_PYRDIN PYR1opt.dmaout final
iteration	iteration
First order induced moments (atomic units) Atom LabelQ10Q1cQ1sN M110.008990743-0.011086254-0.028650718C C12-0.00063404-0.039779257-0.054350889C C23-0.003789783-0.029083768-0.012526122C C340.0012936610.024660883-0.035313287C C25-0.0032138320.016425376-0.047905793C C156-0.0002420760.075838134-0.046018475H H170.006698996-0.011611686-0.003067456H H28-0.001658823-0.015424442-0.003067456H H210-0.0036591130.001953274-0.007078117H H211-0.0037746340.01994962-0.000037930FTrst-order induction energy-2.99010797kJ/mol.Converged induced moments (atomic units)Atom LabelQ1oQ1cQ1sN N110.007457669-0.007754083-0.041300476C C12-0.000886896-0.052311208-0.078088276C C23-0.004847823-0.034178550-0.028088276C C34-0.0006150480.025574195-0.049755066C C25-0.0051084180.024187519-0.00740301H H170.007820067-0.014032519-0.007400301H H28-0.00165530-0.0180953747-0.00230776H H39-0.0055975150.00241838-0.002307676H H3<	First order induced moments (atomic units) Atom Label Q10 Q1c Q1s N N1 1 1 0.014226074 0.023260967 -0.038634199 C c1 2 0.006926994 -0.098149634 -0.075051120 C c2 3 -0.002777219 -0.050396623 -0.043408405 C c3 3 4 -0.00469829 0.006924179 -0.066112972 C c1 5 6 0.003102491 0.079404706 -0.054878888 H H1 1 7 - 0.008676174 -0.02052308 -0.007068704 H H2 2 8 -0.003582663 -0.020731352 -0.011117880 H H3 3 9 -0.004164099 0.002336519 -0.011798518 H H2 4 10 -0.003261779 0.010213061 -0.011702622 H H1 5 11 -0.010716853 0.027068215 0.004596483 First-order induction energy -4.82765443 kJ/mol. Converged induced moments (atomic units) Atom Label Q10 Q1c Q1s N N1 1 1 0.014226075 0.023260967 -0.038634197 C c1 2 3 -0.002777219 -0.050396623 -0.043408404 C c2 3 4 -0.004699828 0.006924179 -0.061537080 C c2 4 5 0.002777219 -0.05396623 -0.043408404 C c2 3 4 -0.004699828 0.006924179 -0.061537080 C c2 4 5 0.002777219 -0.05330623 -0.043408404 C c2 3 4 -0.00469924 0.073352 -0.011178818 H H1 1 7 - 0.00352664 -0.023256487 -0.06812972 C c1 5 6 0.003102491 0.079404706 -0.54878887 H H1 1 7 - 0.00356264 -0.020731352 -0.011178818 H H2 4 10 -0.00356264 -0.020731352 -0.011178818 H H2 4 10 -0.00356179 0.010213061 -0.011702622 H H1 5 11 -0.010716852 0.027068215 0.004596483 Induction energy converged to -4.82765443 kJ/mol with 2 iterations.
24.dampedispersion_PYRDIN fort.12	25.polarizability_PYRDIN fort.12
Initial Lattice Energy: -58.1316 kJ/mol Final Lattice Energy: -61.9302 kJ/mol	Initial Lattice Energy: -61.9760 kJ/mol Final Lattice Energy: -66.7578 kJ/mol
Contributions to lattice energy (eV per primitive unit cell $[kJ/mol]$) (Z = 4) Inter-molecular charge-charge energy= - 0.4502[-10.8588] Total charge-dipole+dipole energy= - 0.0801[-1.9319] Higher multipole interaction energy= - 0.1673[-4.0349] Total isotropic repulsion-dispersion= - 4.0616[-97.9713] Total anisotropic repulsion energy= 2.1917[52.8668] Intermolecular induction energy= 0.0000[0.0000] PV energy.(P = 0.00E+00 Pa)= 0.0000]	Contributions to lattice energy (eV per primitive unit cell $[kJ/mol]$) (Z = 4) Inter-molecular charge-charge energy= - 0.4502[-10.8588] Total charge-dipole+dipole-dipole energy= - 0.0801[-1.9319] Higher multipole interaction energy= - 0.1673[-4.0349] Total isotropic repulsion-dispersion= - 4.0616[-97.9713] Total anisotropic repulsion energy= 2.1917[52.8668] Intermolecular induction energy= - 0.2001[-4.8277] PV energy.(P = 0.00E+00 Pa)= 0.0000[0.0000]

Within the PYRlopt.dmaout file, the induced moments and induction energy are calculate for the initial crystal structure and the final crystal structure. These energies are included in the "Initial Lattice Energy" and "Final Lattice Energy" reports in the fort.12, and hence, as shown above, the lattice energy components differ only in the induction energy between **24.dampedispersion_PYRDIN** and **25.polarisability_PYRDIN**.

References

(1) Price, S. L.; Leslie, M.; Welch, G. W. A.; Habgood, M.; Price, L. S.; Karamertzanis, P. G.; Day, G. M., Modelling Organic Crystal Structures using Distributed Multipole and Polarizability-Based Model Intermolecular Potentials. *Physical Chemistry Chemical Physics* **2010**, 12, (30), 8478-8490.

(2) Day, G. M.; Cooper, T. G.; Cruz-Cabeza, A. J.; Hejczyk, K. E.; Ammon, H. L.; Boerrigter, S. X. M.; Tan, J.; Della Valle, R. G.; Venuti, E.; Jose, J.; Gadre, S. R.; Desiraju, G. R.; Thakur, T. S.; van Eijck, B. P.; Facelli, J. C.; Bazterra, V. E.; Ferraro, M. B.; Hofmann, D. W. M.; Neumann, M.; Leusen, F. J. J.; Kendrick, J.; Price, S. L.; Misquitta, A. J.; Karamertzanis, P. G.; Welch, G. W. A.; Scheraga, H. A.; Arnautova, Y. A.; Schmidt, M. U.; van de Streek, J.; Wolf, A.; Schweizer, B., Significant progress in predicting the crystal structures of small organic molecules - a report on the fourth blind test. *Acta Crystallographica Section B - Structural Science* **2009**, 65, (2), 107-125.

(3) Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.; Robb, M. A.; Cheeseman, J. R.; Montgomery, Jr.; Vreven, T.; Kudin, K. N.; Burant, J. C.; Millam, J. M.; Iyengar, S. S.; Tomasi, J.; Barone, V.; Mennucci, B.; Cossi, M.; Scalmani, G.; Rega, N.; Petersson, G. A.; Nakatsuji, H.; Hada, M.; Ehara, M.; Toyota, K.; Fukuda, R.; Hasegawa, J.; Ishida, M.; Nakajima, T.; Honda, Y.; Kitao, O.; Nakai, H.; Klene, M.; Li, X.; Knox, J. E.; Hratchian, H. P.; Cross, J. B.; Bakken, V.; Adamo, C.; Jaramillo, J.; Gomperts, R.; Stratmann, R. E.; Yazyev, O.; Austin, A. J.; Cammi, R.; Pomelli, C.; Ochterski, J.; Ayala, P. Y.; Morokuma, K.; Voth, G. A.; Salvador, P.; Dannenberg, J. J.; Zakrzewski, V. G.; Dapprich, S.; Daniels, A. D.; Strain, M. C.; Farkas, O.; Malick, D. K.; Rabuck, A. D.; Raghavachari, K.; Foresman, J. B.; Ortiz, J. V.; Cui, Q.; Baboul, A. G.; Clifford, S.; Cioslowski, J.; Stefanov, B. B.; Liu, G.; Liashenko, A.; Piskorz, P.; Komaromi, I.; Martin, R. L.; Fox, D. J.; Keith, T.; Al Laham, M. A.; Peng, C. Y.; Nanayakkara, A.; Challacombe, M.; Gill, P. M. W.; Johnson, B.; Chen, W.; Wong, M. W.; Gonzalez, C.; Pople, J. A. *Gaussian 03*, Gaussian Inc.: Wallingford CT, 2004.

(4) Stone, A. J.; Alderton, M., Distributed multipole analysis - Methods and applications (Reprinted from Molecular Physics, vol 56, pg 1047-1064, 1985). *Molecular Physics* **2002**, 100, (1), 221-233.

(5) Misquitta, A. J.; Welch, G. W. A.; Stone, A. J.; Price, S. L., A first principles solution of the crystal structure of C6Br2ClFH2. *Chemical Physics Letters* **2008**, 456, (1-3), 105-109.

(6) Uzoh, O. G.; Cruz-Cabeza, A. J.; Price, S. L., Is the Fenamate Group a Polymorphophore? Contrasting the Crystal Energy Landscapes of Fenamic and Tolfenamic Acids. *Crystal Growth & Design* 2012, 12, (8), 4230-4239.
(7) Allen, F. H.; Kennard, O.; Watson, D. G., Tables of bond lengths determined by x-ray and neutron diffraction. Part 1. Bond lengths in Organic compounds. *Journal of the Chemical Society-Perkin Transactions 2* 1987, (12), S1-S19.
(8) Bardwell, D. A.; Adjiman, C. S.; Arnautova, Y. A.; Bartashevich, E.; Boerrigter, S. X. M.; Braun, D. E.; Cruz-Cabeza, A. J.; Day, G. M.; Della Valle, R. G.; Desiraju, G. R.; van Eijck, B. P.; Facelli, J. C.; Ferraro, M. B.; Grillo, D.; Habgood, M.; Hofmann, D. W. M.; Hofmann, F.; Jose, K. V. J.; Karamertzanis, P. G.; Kazantsev, A. V.; Kendrick, J.; Kuleshova, L. N.; Leusen, F. J. J.; Maleev, A. V.; Misquitta, A. J.; Mohamed, S.; Needs, R. J.; Neumann, M. A.; Nikylov, D.; Orendt, A. M.; Pal, R.; Pantelides, C. C.; Pickard, C. J.; Price, L. S.; Price, S. L.; Scheraga, H. A.; van de Streek, J.; Thakur, T. S.; Tiwari, S.; Venuti, E.; Zhitkov, I. K., Towards crystal structure prediction of complex organic compounds

- a report on the fifth blind test. Acta Crystallographica Section B-Structural Science 2011, 67, 535-551.
(9) Braun, D. E.; Bhardwaj, R. M.; Florence, A. J.; Tocher, D. A.; Price, S. L., Complex Polymorphic System of Gallic Acid-Five Monohydrates, Three Anhydrates, and over 20 Solvates. Crystal Growth & Design 2013, 13, (1), 19-23.
(10) Kazantsev, A. V.; Karamertzanis, P. G.; Adjiman, C. S.; Pantelides, C. C., Efficient Handling of Molecular Flexibility in Lattice Energy Minimization of Organic Crystals. Journal of Chemical Theory and Computation 2011, 7, (6), 1998-2016.

(11) Wu, H.; Habgood, M.; Parker, J. E.; Reeves-McLaren, N.; Cockcroft, J. K.; Vickers, M.; West, A. R.; Jones, A. G., Crystal structure determination by combined synchrotron powder X-ray diffraction and crystal structure prediction: 1: 1 L-ephedrine D-tartrate. *CrystEngComm* **2013**, 15, (10), 1853-1859.

(12) Day, G. M.; Price, S. L.; Leslie, M., Elastic constant calculations for molecular organic crystals. *Crystal Growth* & *Design* **2001**, 1, (1), 13-27.

(13) Day, G. M.; Price, S. L.; Leslie, M., Atomistic calculations of phonon frequencies and thermodynamic quantities for crystals of rigid organic molecules. *Journal of Physical Chemistry B* **2003**, 107, (39), 10919-10933.
 (14) Day, G. M.; Price, S. L., Properties of Crystalline Organic Molecules. In *Handbook of Elastic Properties of*

Solids, Liquids and Gases. Volume III: Elastic Properties of Solids: Biological and Organic Materials, Earth and Marine Sciences, Levy, M., Ed. Academic Press: New York, 2001; Vol. 3, pp 3-50.

(15) Nyman, J.; Day, G. M., Static and lattice vibrational energy differences between polymorphs. *Crystengcomm* **2015**, 17, (28), 5154-5165.

(16) Nyman, J.; Day, G. M., Modelling temperature-dependent properties of polymorphic organic molecular crystals. *Physical Chemistry Chemical Physics* **2016**, 18, (45), 31132-31143.

(17) Day, G. M.; Price, S. L., A nonempirical anisotropic atom-atom model potential for chlorobenzene crystals. *Journal of the American Chemical Society* **2003**, 125, (52), 16434-16443.

(18) Jongenelis, A. P. J. M.; van der Berg, T. H. M.; Schmidt, J.; van der Avoird, A., Vibron band structure in chlorinated benzene crystals: lattice dynamics calculations and Raman spectra of 1,4-dichlorobenzene. *Journal of Physics: Condensed Matter* **1989**, 1, (31), 5051.

(19) Ridout, J.; Probert, M. R., High-Pressure and Low-Temperature Polymorphism in C-H center dot center dot center dot F-C Hydrogen Bonded Monofluorotoluenes. *Crystal Growth & Design* **2013**, 13, (5), 1943-1948.

(20) Kazantsev, A. V.; Karamertzanis, P. G.; Adjiman, C. S.; Pantelides, C. C.; Price, S. L.; Galek, P. T. A.; Day, G. M.; Cruz-Cabeza, A. J., Successful prediction of a model pharmaceutical in the fifth blind test of crystal structure prediction. *International Journal of Pharmaceutics* **2011**, 418, (2), 168-178.

(21) Misquitta, A. J.; Stone, A. J. *CamCASP: a program for studying intermolecular interactions and for the calculation of molecular properties in distributed form*, University of Cambridge, <u>http://www-stone.ch.cam.ac.uk/programs.html#CamCASP:</u> 2007.

(22) Braun, D. E.; Ardid-Candel, M.; D'Oria, E.; Karamertzanis, P. G.; Arlin, J. B.; Florence, A. J.; Jones, A. G.; Price, S. L., Racemic Naproxen: A Multidisciplinary Structural and Thermodynamic Comparison with the Enantiopure Form. *Crystal Growth & Design* **2011**, 11, (12), 5659-5669.

(23) Aina, A. A.; Misquitta, A. J.; Price, S. L., From dimers to the solid-state: Distributed intermolecular force-fields for pyridine. *The Journal of Chemical Physics* **2017**, 147, (16), 161722.

(24) Misquitta, A. J.; Stone, A. J., Ab Initio Atom-Atom Potentials Using CAMCASP: Theory and Application to Many-Body Models for the Pyridine Dimer. *Journal of Chemical Theory and Computation* 2016, 12, (9), 4184-4208.
(25) Aina, A.; Misquitta, A.; Price, S., From dimers to the solid-state: Distributed intermolecular force-fields for pyridine. *Journal of Chemical Physics* 2017, 147, (16).